This document introduces the implementation of the Armstrong-Frederick Steel Model.

### 1 Yield Function

A von Mises yielding function is used.

$$F = \sqrt{\frac{3}{2}} \left| \eta \right| - k,\tag{1}$$

in which  $\eta = s - \beta$  is the shifted stress, *s* is the stress deviator,  $\beta$  is the back stress and *k* is the isotropic hardening stress.

#### 2 Flow Rule

The associated plasticity flow is adopted. The plastic strain rate is then

$$d\varepsilon^p = \gamma \frac{\partial F}{\partial \sigma} = \sqrt{\frac{3}{2}} \gamma n, \tag{2}$$

where  $n = \frac{\eta}{|\eta|}$ . The corresponding accumulated plastic strain rate is

$$\mathrm{d}p = \sqrt{\frac{2}{3}}\mathrm{d}\varepsilon^p : \mathrm{d}\varepsilon^p = \gamma. \tag{3}$$

## 3 Hardening Rules

An exponential function with a linear component is used for isotropic hardening stress.

$$k = \sigma_y + k_1 p + k_s - k_s e^{-mp}.$$

The corresponding derivative is

$$\frac{\mathrm{d}k}{\mathrm{d}\gamma} = k_l + k_s m e^{-mp}.\tag{5}$$

The rate form of back stress  $\beta = \sum \beta^i$  is defined as

$$\mathrm{d}eta^i = \sqrt{rac{2}{3}}a^i\mathrm{d}arepsilon^p - b^ieta^i\mathrm{d}p.$$

In terms of  $\gamma$ , it is  $d\beta^i = a^i \gamma n - b^i \gamma \beta^i$ . The incremental form is thus

$$\beta^{i} = \beta^{i}_{n} + a^{i}\gamma n - b^{i}\gamma\beta^{i}, \qquad \beta^{i} = \frac{\beta^{i}_{n} + a^{i}\gamma n}{1 + b^{i}\gamma}.$$
(6)

#### 4 Incremental Form

The shifted stress can be computed as

$$\eta = s - \beta = 2GI_d \left(\varepsilon^{tr} - \varepsilon_n^p - \sqrt{\frac{3}{2}}\gamma n\right) - \beta = s^{tr} - \sqrt{6}G\gamma n - \sum \frac{\beta_n^i + a^i\gamma n}{1 + b^i\gamma}$$
(7)

with  $s^{tr} = 2GI_d \left(\varepsilon^{tr} - \varepsilon_n^p\right)$ . Hence,

$$\left|\eta\right|n + \sqrt{6}G\gamma n + \sum \frac{a^{i}\gamma}{1+b^{i}\gamma}n = s^{tr} - \sum \frac{\beta_{n}^{i}}{1+b^{i}\gamma}, \qquad \left|\eta\right| + \sqrt{6}G\gamma + \sum \frac{a^{i}\gamma}{1+b^{i}\gamma} = \left|s^{tr} - \sum \frac{\beta_{n}^{i}}{1+b^{i}\gamma}\right|.$$

Eventually,

$$\left|\eta\right| = \left|s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}\right| - \sqrt{6}G\gamma - \sum \frac{a^i \gamma}{1 + b^i \gamma}.$$
(8)

Then  $\eta$  can be expressed as,

$$\eta = \frac{\left|s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}\right| - \sqrt{6}G\gamma - \sum \frac{a^i \gamma}{1 + b^i \gamma}}{\left|s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}\right|} \left(s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}\right).$$

It is equivalent to

$$\eta = \left( \left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right| - \sqrt{6}G\gamma - \sum \frac{a^i \gamma}{1 + b^i \gamma} \right) u.$$

It is easy to see that n = u with  $u = \frac{s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}}{\left|s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}\right|}$ . The derivatives of u are

$$\frac{\partial u}{\partial \gamma} = \frac{1}{\left|s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}\right|} \sum \frac{b^i}{(1 + b^i \gamma)^2} \left(\beta_n^i - u : \beta_n^i u\right), \qquad \frac{\partial u}{\partial \varepsilon^{tr}} = \frac{1}{\left|s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}\right|} 2G\left(I_d - u \otimes u\right).$$
(9)

# 5 Scalar Equation Iteration

The yield function is then

$$F = \sqrt{\frac{3}{2}} \left( \left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right| - \sqrt{6}G\gamma - \sum \frac{a^i \gamma}{1 + b^i \gamma} \right) - \left( \sigma_y + k_l \left( p_n + \gamma \right) + k_s \left( 1 - e^{-m(p_n + \gamma)} \right) \right).$$
(10)

The corresponding derivative is then

$$\frac{\partial F}{\partial \gamma} = \sqrt{\frac{3}{2}} \sum \frac{b^i u : \beta_n^i - a^i}{(1 + b^i \gamma)^2} - 3G - k_l - k_s m e^{-m(p_n + \gamma)}, \qquad \frac{\partial F}{\partial \varepsilon^{tr}} = \sqrt{6} G u : I_d = \sqrt{6} G u. \tag{11}$$

Note  $u : I_d = u$  since u is a deviator.

# 6 Consistent Tangent Stiffness

For stiffness,  $\varepsilon^{tr}$  is now varying, then

$$\frac{\partial F}{\partial \varepsilon^{tr}} + \frac{\partial F}{\partial \gamma} \frac{d\gamma}{d\varepsilon^{tr}} = 0, \qquad \frac{d\gamma}{d\varepsilon^{tr}} = -\left(\frac{\partial F}{\partial \gamma}\right)^{-1} \frac{\partial F}{\partial \varepsilon^{tr}}.$$
(12)

Since the stress can be written as

$$\sigma = E(\varepsilon^{tr} - \varepsilon^p) = E(\varepsilon^{tr} - \varepsilon^p_n - \Delta\varepsilon^p) = E(\varepsilon^{tr} - \varepsilon^p_n) - \sqrt{6}G\gamma u.$$
(13)

The derivative is

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\varepsilon^{tr}} = E - \sqrt{6}G\left(\gamma\frac{\partial u}{\partial\varepsilon^{tr}} + \left(u + \gamma\frac{\partial u}{\partial\gamma}\right)\frac{\mathrm{d}\gamma}{\mathrm{d}\varepsilon^{tr}}\right) \\
= E - \sqrt{6}G\gamma\frac{\partial u}{\partial\varepsilon^{tr}} + 6G^2\left(u + \gamma\frac{\partial u}{\partial\gamma}\right)\left(\frac{\partial F}{\partial\gamma}\right)^{-1}u.$$
(14)